

that the cations depart significantly from S_4 symmetry. Zaslow and Rundle²⁰ had commented that the versatility of form of the tetraphenylarsonium cation helps to account for its effectiveness as a precipitating agent. Our results serve to reinforce this statement and help to explain the effectiveness of this large cation in precipitating trinuclear halogeno anions of Re(III) which have fewer than 12 halogen atoms.¹⁵

Ionic Packing.—The anions are well shielded from one another by the bulky cations as can be seen from Figure 1 and there are no interionic Cl---Cl contacts within 4.5 Å. The closest Cl---C contacts made by individual chlorine atoms range from 3.2 Å for Cl(10) to 3.8 Å for Cl(7), which, considering the large errors

in carbon positions, are reasonable van der Waals distances.

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The Crystal Structure of Bis(cesium) Tri- μ -bromo-octabromotrirhenate(III), $\text{Cs}_2\text{Re}_3\text{Br}_{11}$

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A three-dimensional crystal structure analysis of the compound of composition $\text{Cs}_2\text{Re}_3\text{Br}_{11}$ establishes that it contains the trinuclear anion $(\text{Re}_3\text{Br}_{11})^{2-}$ which possesses the same general structural features as other rhenium(III) species of the general formula $(\text{Re}_3\text{X}_{9+n})^{n-}$ where X is Br or Cl. As with the $(\text{Re}_3\text{Cl}_{11})^{2-}$ anion, Re-Re bond lengths are unequal, those involving the deficient Re atom being 2.43 Å and the remaining Re-Re bond 2.49 Å. A new feature is that the bonds which the deficient Re atom makes to its off-plane bromine atoms are greatly shortened to 2.38 Å and are bent toward one another to make an angle of 133°.

Introduction

The existence of trinuclear Re(III) bromide species analogous to the chloride compounds already known¹⁻⁴ was first reported by Fergusson and Robinson⁵ and independently by Cotton and Lippard.⁶ The preparative and spectral properties of some of these compounds have been discussed by Robinson and Fergusson⁷ and by Cotton, Lippard, and Mague.⁸ It is as a part of our systematic studies of the structures of trinuclear Re(III) species that we have determined in detail the crystal structure of $\text{Cs}_2\text{Re}_3\text{Br}_{11}$, a preliminary account of which has already been given.⁹ The crystal structure of one other compound containing the neutral Re_3Br_9 species has been reported by Cotton and Lippard.¹⁰

(1) W. T. Robinson, J. E. Fergusson, and B. R. Penfold, *Proc. Chem. Soc.*, 116 (1963); J. A. Bertrand, F. A. Cotton, and W. A. Dollase, *J. Am. Chem. Soc.*, **85**, 1349 (1963); *Inorg. Chem.*, **2**, 1166 (1963).

(2) J. E. Fergusson, B. R. Penfold, and W. T. Robinson, *Nature*, **201**, 181 (1964).

(3) F. A. Cotton and J. T. Mague, *Inorg. Chem.*, **3**, 1094 (1964).

(4) F. A. Cotton and J. T. Mague, *ibid.*, **3**, 1402 (1964).

(5) J. E. Fergusson and B. H. Robinson, *Proc. Chem. Soc.*, 189 (1964).

(6) F. A. Cotton and S. J. Lippard, *J. Am. Chem. Soc.*, **86**, 4497 (1964).

(7) B. H. Robinson and J. E. Fergusson, *J. Chem. Soc.*, 5683 (1964, Suppl. 1).

(8) F. A. Cotton, S. J. Lippard, and J. T. Mague, *Inorg. Chem.*, **4**, 508 (1965).

(9) M. Elder and B. R. Penfold, *Nature*, **205**, 276 (1965).

(10) F. A. Cotton and S. J. Lippard, *Inorg. Chem.*, **4**, 59 (1965).

Experimental Section

Small black needle-shaped crystals were prepared following the method of Fergusson and Robinson.⁵ The crystals were completely stable in air.

Crystal Data.— $\text{Cs}_2\text{Re}_3\text{Br}_{11}$, formula weight 1704, orthorhombic, space group $Pnma$ from systematic absences of X-ray reflections; $a = 9.53 \pm 0.01$ Å, $b = 16.114 \pm 0.007$ Å, $c = 13.744 \pm 0.007$ Å ($\lambda = 1.5405, 1.5443$ for $\text{Cu K}\alpha_{1,2}$ radiation), $U = 2109.5$ Å³, $D_x = 5.38$ g cm⁻³, $Z = 4$, symmetry of anion $C_s (=m)$. Linear absorption coefficients are 865 cm⁻¹ for $\text{Cu K}\alpha$ and 459 cm⁻¹ for $\text{Mo K}\alpha$ radiation.

Reciprocal lattice levels ($0-6kl$) were recorded by the equi-inclination Weissenberg method using $\text{Cu K}\alpha$ radiation, and levels ($h\bar{k}0-1$), ($h0l$), and ($h\bar{k}k-h\bar{k}(k+4)$) were recorded on the Buerger precession camera using $\text{Mo K}\alpha$ radiation. Intensities were estimated visually by comparison with a suitably calibrated scale. The Weissenberg data totalled 1713 independent reflections of which 323 were recorded as unobserved and given an estimated intensity one-third of the minimum observable intensity. These Weissenberg data were used for the structure refinement, supplemented by 178 precession reflections with $h > 6$, of which 30 were recorded as unobserved. The two crystals used were each mounted on a glass fiber and showed no signs of deterioration after prolonged exposure to the atmosphere and to X radiation. They were of a very similar size, needles of length about 0.5 mm and rectangular cross section 0.095 mm \times 0.070 mm. Corrections for absorption were applied assuming a cylindrical crystal of radius 0.04 mm ($\mu R = 3.5$) for the Weissenberg data and assuming a spherical crystal of the same radius ($\mu R = 1.8$) for the precession data.

Structure Determination

A three-dimensional trial structure was based upon a self-consistent set of atomic positions obtained from solutions of the Patterson function projected down each of the principal axes. Because the rhenium triangle is nearly parallel to (101) there is good resolution of atoms in the projections on (001) and (100). The cesium atom, one rhenium atom, and four bromine atoms were located on general (8-fold) positions, and the one remaining rhenium atom and three remaining bromine atoms on mirror planes normal to *b*, positions of 4-fold multiplicity. Four successive three-dimensional difference maps and subsequent sets of structure factors were computed, in the course of which all atoms were more precisely located and the *R* index dropped from 0.40 to 0.27 for all Weissenberg data. Next, five cycles of least-squares refinement were performed using the same data and *R* was further reduced to 0.17. The block diagonal approximation was used with a two-parameter weighting scheme suggested by Mills and Rollett.¹¹ The weight is expressed as

$$w = 1 / \left\{ 1 + \left(\frac{F_o - b}{a} \right)^2 \right\}$$

with *a* = 90 and *b* = 60 in this case. Atomic scattering factors used were, for Cs and Re, from the Thomas-Fermi-Dirac statistical model and, for Br, from the self-consistent field model, all corrected for anomalous dispersion,¹² the actual values used being $f = \sqrt{\{f_o + \Delta f''\}^2 + (\Delta f')^2}$.

At this stage in the refinement all data were scaled to $\Sigma F_o = \Sigma |F_c|$ for individual layers. The precession data covering *h* = 7-11 were added, and all unobserved reflections were removed. The remaining 1538 reflections were used in the computation of a final difference map. This showed no evidence of anisotropy for any atoms, so all atoms were treated as isotropic in the final cycles of least squares. Three cycles were run and the *R* index fell from 0.18 to its final value of 0.145.

Final atomic coordinates and isotropic thermal parameters are listed in Table I. The *B* values have no absolute significance since absorption corrections calculated on the basis of a regular cylinder can only be approximate. Corresponding structure factors are given in Table II for all observed reflections. Interatomic distances and bond angles involving the anions are listed in Table III.

Appended to Table III are estimated mean standard deviations based on 1.5 times the indications of the least-squares refinement. Although somewhat arbitrary, this figure is reasonable on the basis of past experience with the block diagonal approximation in similar structure analyses and also because of the internal consistency of chemically equivalent bonds and angles, which will be expected to agree within $3\sqrt{2}\sigma$

(11) O. S. Mills and J. S. Rollett, "Computing Methods and the Phase Problem in X-ray Crystal Analysis," Pergamon Press, London, 1961, p 117.

(12) "International Tables for X-ray Crystallography," Vol. III, the Kynoch Press, Birmingham, England, 1962, pp 206-216.

TABLE I
ATOMIC PARAMETERS FOR THE ASYMMETRIC
UNIT OF Cs₂Re₃Br₁₁^a

	<i>x</i>	<i>y</i>	<i>z</i>	<i>B</i> , Å ²
Re(1)	0.3020	0.1725	0.4669	2.3
Re(2)	0.4461	0.2500	0.5814	2.5
Cs	0.3368	0.0697	0.1202	4.2
Br(1)	0.0961	0.1478	0.5714	3.4
Br(2)	0.4869	0.1462	0.3427	2.8
Br(3)	0.2057	0.0304	0.3832	3.6
Br(4)	0.3354	0.2500	0.7355	4.2
Br(5)	0.6864	0.2500	0.5299	3.9
Br(6)	0.4572	0.0915	0.5863	3.9
Br(7)	0.1538	0.2500	0.3452	3.2

^a Atom numbering corresponds to that of Figure 1.

(=4.2σ). The agreement is well within this figure for all sets except for the two pairs of angles Br(1)-Re(1)Br(3) and Br(2)Re(1)Br(3) (difference of 4.2σ), and Br(1)Re(1)Br(6) and Br(2)Re(1)Br(6) (difference of 7σ). Intermolecular forces could well influence these angles, Br(2) being involved in the closest of all intermolecular contacts.

Discussion

Crystal Structure.—The crystal structure consists of discrete (Re₃Br₁₁)²⁻ anions and Cs⁺ cations. A view of the structure projected down the *c* axis is shown in Figure 1. The anions occur on mirror planes and are related by the two glide planes. The cesium atom has nine bromine neighbors at less than 4.0 Å, five at 3.57-3.66 Å, and a further four at 3.82-3.88 Å. There is no discernible pattern in the arrangement of the bromine atoms around the cesium atom, but this is predictable since the anion packing must determine the structure. None of the distances is abnormally short compared with the value in CsBr (3.72 Å), the shortest of 3.57 Å being comparable with the shortest Cs---Cl distance of 3.3 Å observed in Cs₃Re₃Cl₁₂.¹ The two shortest intermolecular Br---Br contacts are indicated in Table III. These values lie between the van der Waals distance of 3.9 Å and those for the short intramolecular nonbonded contacts, of which the shortest is 3.29 Å.

The (Re₃Br₁₁)²⁻ Anion.—The anion is based on the now familiar structural framework of the (Re₃Cl_{9+n})ⁿ⁻ series¹⁻⁴ and the Re₃Br₉ species.¹⁰ It closely resembles the (Re₃Cl₁₁)²⁻ anion,² possessing only two out of the three possible terminal bromine atoms in the plane of the Re₃ triangle. Thus, Re(2) is bonded to only four bromine atoms compared with the five for Re(1) and is termed "deficient."

The asymmetric unit includes one-half of the anion, which lies on a crystallographic mirror plane passing through the unique rhenium atom, Re(2), its two terminal bromine atoms, Br(4) and Br(5), and the bromine, Br(7), that bridges Re(1) and Re(1'). There are no significant deviations from the least-squares plane (calculated by the method of Schomaker, *et al.*)¹³ for the atoms Re(1), Re(2), Br(3), Br(6), Br(7) of

(13) V. Schomaker, J. Waser, R. E. Marsh, and G. Bergman, *Acta Cryst.* **12**, 600 (1959).

TABLE III
INTERATOMIC DISTANCES AND BOND ANGLES
INVOLVING THE $(\text{Re}_3\text{Br}_{11})^{2-}$ ANIONS^a

Bond	Length, A		Value, deg
		Angle	
Re(1)-Re(1')	2.496	Re(1)Br(6)Re(2)	56.7
Re(1)-Re(2)	2.433	Re(1)Re(2)Br(6)	61.8
Re(1)-Br(1)	2.46	Re(2)Re(1)Br(6)	61.5
Re(1)-Br(2)	2.49	Re(1)Re(2)Re(1')	61.7
Re(1)-Br(3)	2.72	Re(1)Re(1')Re(2)	59.1
Re(2)-Br(4)	2.37	Re(1)Br(7)Re(1')	59.4
Re(2)-Br(5)	2.40	Re(1)Re(1')Br(7)	60.3
Re(1)-Br(6)	2.57	Br(1)Re(1)Br(7)	91.2
Re(1)-Br(7)	2.52	Br(1)Re(1)Br(3)	80.9
Re(2)-Br(6)	2.56	Br(1)Re(1)Br(6)	90.2
Br(1)---Br(1')	3.29	Br(2)Re(1)Br(7)	91.5
Br(2)---Br(2')	3.34	Br(2)Re(1)Br(3)	78.8
Br(1)---Br(4)	3.73	Br(2)Re(1)Br(6)	86.8
Br(2)---Br(5)	3.76	Br(4)Re(2)Br(5)	133.7
Br(1)---Br(3)	3.37	Br(4)Re(2)Br(6)	89.7
Br(2)---Br(3)	3.31	Br(5)Re(2)Br(6)	88.2
Br(5)---Br(6)	3.45		
Br(4)---Br(6)	3.47		
Br(2)---Br(6)	3.48		
Br(2)---Br(7)	3.59		
Br(1)---Br(6)	3.56		
Br(1)---Br(7)	3.56		
Br(2)---Br(7) ^b	3.46		
Br(4)---Br(5) ^b	3.52		

^a Estimated standard deviations based upon 1.5 times the indication of the least squares refinement are: Re(1)-Re(2), $\sigma = 0.008$ A; Re(1)-Re(1'), $\sigma = 0.011$ A; Re-Br, $\sigma = 0.012$ A; Br---Br, $\sigma = 0.016$ A; all angles, $\sigma \leq 0.5^\circ$. ^b Denotes intermolecular contact between bromines of adjacent anions lying on the same mirror plane.

The most interesting dimensional features are the nonequality of the Re-Re bond lengths and the length and angular relation of the terminal Re-Br bonds to the deficient rhenium. The Re-Re bond lengths parallel those in $(\text{Re}_3\text{Cl}_{11})^{2-}$. The two short bonds, associated with the deficient rhenium atom, are 2.433 A (2.435 A in $(\text{Re}_3\text{Cl}_{11})^{2-}$). The bond between the nondeficient rhenium atoms is 2.496 A (2.483 A in $(\text{Re}_3\text{Cl}_{11})^{2-}$). (In $(\text{Re}_3\text{Br}_{12})^{3-14}$ and Re_3Br_9 , where the Re_3 triangles are equilateral, the corresponding mean lengths are 2.49 and 2.465 A, respectively.) The observation in $(\text{Re}_3\text{Cl}_{11})^{2-}$ of an increase in Re-Re bond strength associated with a rhenium atom that forms only six σ bonds instead of seven is thus strikingly confirmed. This increased bond strength has been accounted for in terms of greater directional character of the changed σ -bond framework.¹⁵

In all the trimeric rhenium compounds studied so far, the angle subtended by a pair of off-plane halogen atoms at a nondeficient rhenium atom has been within the range 155-160°. Representative values are 158° for $(\text{Re}_3\text{Cl}_{12})^{3-}$, 155° for $(\text{Re}_3\text{Br}_{12})^{3-}$, and 159° for $(\text{Re}_3\text{Br}_{11})^{2-}$ and $(\text{Re}_3\text{Cl}_{11})^{2-}$. When the same angle is considered for halogen atoms bonded to a deficient rhenium atom, it is found that in $(\text{Re}_3\text{Cl}_{11})^{2-}$ and in Re_3Br_9 the values (153 and 155°, respectively) are close to the above figures. (In $(\text{Re}_3\text{Cl}_{11})^{2-}$ the deficient Re atom is probably associated with an in-plane water molecule.²) In $(\text{Re}_3\text{Br}_{11})^{2-}$, however, the angle Br-(4)Re(2)Br(5) is only 133°. This reduction does not

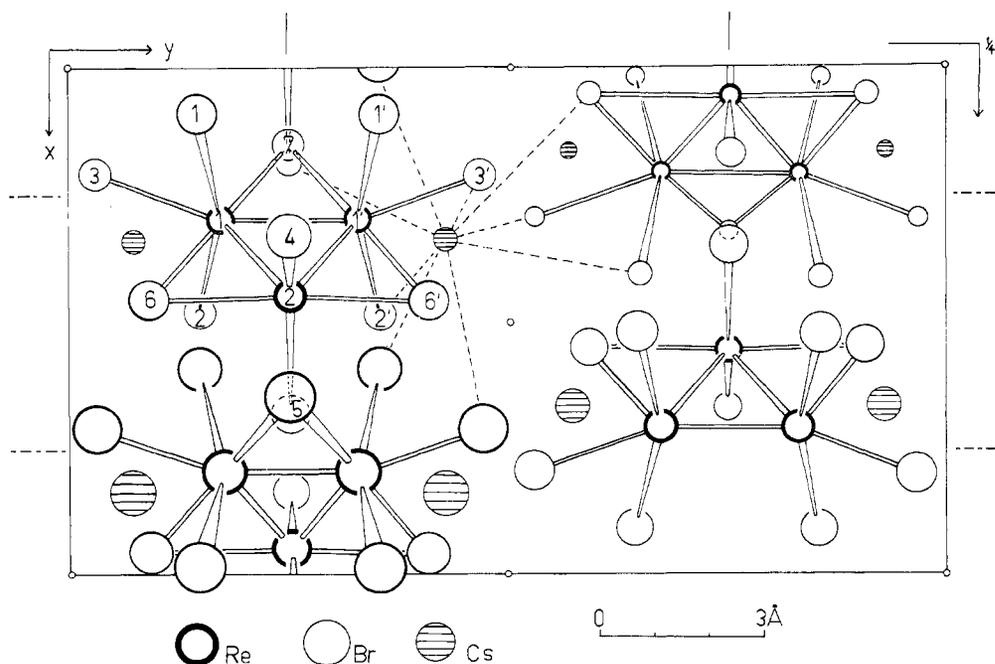


Figure 1.—The $\text{Cs}_2\text{Re}_3\text{Br}_{11}$ structure as viewed down the c axis. Broken lines indicate Cs---Br contacts, some of which are directed to anions related by translations $\pm c$ to those shown.

one anion. The equation of the plane is $-0.7579X - 0.0118Y + 0.6523Z + 0.0019 = 0$, where coordinates are expressed in angstrom units. Within limits of error this plane is a plane of symmetry for the anion, which thus has effective C_{2v} symmetry.

appear to have been brought about by intermolecular forces. The only intermolecular distance less than 4.0

(14) G. J. Gainsford, B. R. Penfold, and M. Elder, to be published.

(15) J. E. Fergusson, B. R. Penfold, M. Elder, and B. H. Robinson, *J. Chem. Soc.*, 5500 (1965).

A involving Br(4) or Br(5) is the Br(4)---Br(5) distance of 3.52 Å between ions related by the *a* glide plane. This contact and the other short intermolecular distance are dependent upon the inclination of the Re_3 triangle to the (001) plane, but are not necessarily related to the size of the Br(4)Re(2)Br(5) angle. When the short intramolecular contacts are considered it is clear that a marked reduction in repulsive forces has been achieved by the closing of this angle. First we note that the distances Br(4)---Br(6) and Br(5)---Br(6) are virtually independent of the angle under consideration. This leaves the contacts Br(1)---Br(4) and Br(2)---Br(5) as the important factors. These two independent distances average 3.74 Å, which, when compared with the Br(1)---Br(1') distance of 3.31 Å and the corresponding 3.38 Å in Re_3Br_9 , indicates a considerable decrease in bromine-bromine repulsive forces across the face of the anion.

Accompanying the decrease in the Br(4)Re(2)Br(5) angle is a corresponding shortening of the two off-plane terminal bonds to a mean value of 2.38 Å, compared with the mean value of 2.47 Å for the bonds to Re(1). This difference of 0.09 Å is about 7σ and therefore highly significant. It may be accounted for by

postulating some π bonding between filled 3d orbitals of the terminal bromines and the rhenium 5p or 5s5p hybrid lying in the plane of the Re_3 triangle and unused in the absence of an in-plane terminal bromine. It is significant that shortening of the terminal off-plane bonds has not occurred in the isolated Re_3Br_9 ¹⁰ species where *all* rhenium atoms are "deficient." Neither has the angle subtended by the pair of off-plane bromine atoms decreased. Additional π bonding cannot therefore be postulated in this case. It is hard to believe that intermolecular packing forces (which appear to be quite weak in the complex in which Re_3Br_9 occurs) could cause such wide variations as exist between the dimensions of $(Re_3Br_{11})^{2-}$ and neutral Re_3Br_9 . For a full understanding of these variations we must await a quantitative bonding theory which considers the trinuclear cluster as a whole.

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CONTRIBUTION FROM THE LAWRENCE RADIATION LABORATORY AND DEPARTMENT OF CHEMISTRY,
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The Crystal Structure of Chlorothiodiazyl Chloride, $S_3N_2Cl_2$ ¹

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$S_3N_2Cl_2$ crystallizes in the monoclinic space group $P2_1$ with cell dimensions $a = 6.546$ Å, $b = 8.600$ Å, $c = 5.508$ Å, and $\beta = 102.37^\circ$; there are two formula units per unit cell. The material is a salt consisting of a chloride anion and a $S-N-S-N-S-$

Cl^+ cation. The sulfur and nitrogen atoms form a puckered five-membered ring.

Introduction

$S_3N_2Cl_2$ was first identified by Demarcay,² who prepared it by allowing SCl_2 or S_2Cl_2 to react with S_4N_4 . Meuwesen³ prepared $S_3N_2Cl_2$ by allowing $(NSCl)_3$ to react with S_2Cl_2 . More recently Jolly, Maguire, and Rabinovich⁴ produced the material by refluxing a suspension of ammonium chloride in S_2Cl_2 . An X-ray crystal structure analysis on a crystal prepared by Jolly's method was undertaken to determine the molecular structure of this compound.

Experimental Section

$S_3N_2Cl_2$ is a yellow crystalline salt which decomposes in the air. The salt was removed from its container in a nitrogen-

filled drybox, crushed, and sieved. Several crystal fragments of about 0.3 mm were introduced into tapered quartz capillaries of about the same size and sealed. With patient use of a microscope and a file, a single crystal fragment was isolated and wedged in the capillary; the capillary was fire sealed to a final length of about 1 cm. The crystal fragment was an irregular plate with dimensions approximately $0.3 \times 0.4 \times 0.2$ mm. Some decomposition of the crystal surface did occur as the crystal lost its yellow color and took on a blackish tinge; however, the crystal diffracted X-rays well and no further decomposition was encountered. The crystal was oriented with the monoclinic *b* axis parallel to the length of the capillary.

A General Electric XRD-5 apparatus equipped with a quarter-circle Eulerian cradle type goniostat, a scintillation counter, a pulse height discriminator, and a molybdenum X-ray tube was used to measure the cell dimensions and the intensities. The X-ray tube was operated at 20 ma and 40 kv. The primary X-ray beam was filtered with a Zr filter.

The measured monoclinic cell dimensions based on $\lambda_{K\alpha_1} = 0.70926$ Å are: $a = 6.546 \pm 0.007$ Å, $b = 8.600 \pm 0.004$ Å, $c = 5.508 \pm 0.005$ Å, and $\beta = 102.37 \pm 0.03^\circ$. The observed density was 2.0 g/ml, suggesting two formula units per unit cell;

(1) Work done under the auspices of the U. S. Atomic Energy Commission.

(2) E. Demarcay, *Compt. Rend.*, **92**, 726 (1881).

(3) A. Meuwesen, *Ber.*, **65**, 1731 (1932).

(4) W. L. Jolly, K. D. Maguire, and D. Rabinovich, *Inorg. Chem.*, **2**, 1304 (1963).